

Non-trivial analytical solutions in 2- and 3-dimensional scalar field models

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Here we demonstrate in explicit form the analytical solutions (instantons) for the scalar field models in $S^1 \otimes R^1$ space-time and analytical solutions (solitons) for the scalar field models in $S^1 \otimes R^1 \otimes R^1$ space-time.

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Let us consider 2-dimensional scalar model in $S^1 \otimes R^1$ space-time ($-\frac{L}{2} \leq x \leq \frac{L}{2}$, $-\infty < t < \infty$):

$$L = \frac{1}{2} \partial_\mu \varphi \partial_\mu \varphi - V(\varphi), \quad \mu = 0, 1, \quad (1)$$

$$V(\varphi) = \lambda(1 - \cos(\rho\varphi)), \quad (1.1)$$

$$V(\varphi) = \lambda(\rho^2 - \varphi^2)^2, \quad (1.2)$$

where $\varphi(x, t)$, λ , ρ are real.

This models admit instantons (non-trivial classical solutions of the field equations in imaginary time $\tau = it$) [1]:

$$\varphi^{inst}(x, \tau) = \pm \frac{4}{\rho} \arctan e^{(\tau - \tau_0)\rho\sqrt{\lambda}}, \quad (2.1)$$

$$\varphi^{inst}(x, \tau) = \pm \tanh\left(\rho\sqrt{2\lambda}(\tau - \tau_0)\right), \quad (2.2)$$

where τ_0 is arbitrary parameter. These space-homogeneous solutions can be used for the description of the tunneling transitions between neighbor classically degenerated vacua of $V(\varphi)$ [2].

The (Euclidean) actions on instantons can be easy calculated [1]:

$$S_I[\varphi^{inst}(x, \tau)] = \frac{8\sqrt{\lambda}}{\rho} L, \quad (3.1)$$

$$S_I[\varphi^{inst}(x, \tau)] = \frac{4\sqrt{2\lambda}}{3} \rho^3 L. \quad (3.2)$$

Let us consider 3-dimensional scalar model in $S^1 \otimes R^1 \otimes R^1$ space-time ($-\frac{L}{2} \leq x \leq \frac{L}{2}$, $-\infty < y < \infty$, $-\infty < t < \infty$):

$$L = \frac{1}{2} \partial_\mu \varphi \partial_\mu \varphi - V(\varphi), \quad \mu = 0, 1, 2, \quad (4)$$

where $V(\varphi)$ are defined in (1.1) and (1.2).

It is well-known, that instanton solutions in D dimensions formally coincide with static solitons in $D + 1$ dimensions. Therefore we can immediately write exact solitons solutions:

$$\varphi^{sol}(x, y, t) = \pm \frac{4}{\rho} \arctan e^{(y - y_0)\rho\sqrt{\lambda}}, \quad (5.1)$$

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$$\varphi^{sol}(x, y, t) = \pm \tanh \left(\rho \sqrt{2\lambda} (y - y_0) \right). \quad (5.2)$$

The energies of solitons (5.1) and (5.2) formally coincide with (Euclidean) actions on instantons (3.1) and (3.2) correspondingly:

$$E[\varphi^{sol}(x, y, t)] = \frac{8\sqrt{\lambda}}{\rho} L, \quad (6.1)$$

$$E[\varphi^{sol}(x, y, t)] = \frac{4\sqrt{2\lambda}}{3} \rho^3 L. \quad (6.2)$$

Depending on time solitons can be obtained from solutions (5.1) and (5.2) by transition to another frame of reference:

$$\varphi^{sol}(x, y, t) = \pm \frac{4}{\rho} \arctan \left[\exp \left(\frac{(y - y_0) - ut}{\sqrt{1 - u^2}} \rho \sqrt{\lambda} \right) \right], \quad (7.1)$$

$$\varphi^{sol}(x, y, t) = \pm \tanh \left[\frac{(y - y_0) - ut}{\sqrt{1 - u^2}} \rho \sqrt{2\lambda} \right], \quad (7.2)$$

where u is a velocity of new frame of reference.

References

- [1] R.G.Shulyakovsky. *In Proceed. of the 12th Annual Seminar "Nonlinear Phenomena in Complex Systems"* (May 17-20, 2005, Minsk, Belarus), 169.
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